



A Review Research on Rotational Symmetries with Magnetic Domains

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ABSTRACT

Due to their numerous industrial applications, magnetic films have attracted a lot of attention for decades. Their domain structure is increasingly important to their present functionality. Due to the conflict between short-range attractive and long-range repulsive interactions, magnetic films typically develop complicated domain patterns with distinctive structures at different length scales. Despite having orientation spin order in each of them, the group of domains is topologically disordered. Finding new ordering from the ostensibly disordered structures would be an interesting topic because simplifying complexity is the essential step to understanding and changing nature. It is well known that scattering techniques are effective instruments for finding ordering.

KEYWORDS: Magnetic films, dimensions, rotational systems.

INTRODUCTION

In nature, patterns frequently arise through self-organization in biological, chemical, and physical systems [1]. Scientists from a wide range of fields have focused intense theoretical and experimental research efforts on it because of the similarities in the fundamental mechanisms by which nonlinearities combine to generate spatial or temporal patterns. Since a large deal of curiosity was generated by phase transition phenomena, which are typically accompanied by complex patterns emerging from straightforward uniform mixtures, it has been the mainstay of non-equilibrium physics for decades. Dynamics is one of the key characteristics connected to non-equilibrium [2, 3]. The process of pattern generation must be controlled by some nonlinear mechanism that shifts the system to a new state, frequently exhibiting a spatial or temporal pattern, when some driving force disturbs a spatially uniform system in a stable state, close to the commencement of instability. Phenomenological Ginzburg-Landau-type theories can be used to explain weakly nonlinear patterns in a well-developed theoretical framework. Under various conditions, magnetic materials exhibit a variety of intricate patterns and solitary structures over a wide range of length scales [4]. For instance, different structures can be seen in magnetic fluids when exposed to external magnetic fields [5, 6]. Remarkable one-dimensional patterns can be created when the fields are oriented parallel to the plane of the films, and two-dimensional lattices can be created when the fields are oriented perpendicular to the plane of the films [7]. In the condition described above, the pattern development can be thought of as a dissipative dynamical process [8, 9]. Since complexity in the structure is frequently associated with sophisticated functionalities, novel magnetic devices, such as magnetic multilayers, exhibit surprising structural and magnetic properties even in the context of mechanical equilibrium patterns. These devices result from a system whose free energy has a complexed structure with many local minima corresponding to metastable states. Instead of the condition of least energy, metastable states are where hysteresis phenomena originate [10]. The non-equilibrium thermodynamics field, which still has many open questions, is where the mechanisms behind hysteresis reside. However, a straightforward qualitative explanation can be offered. When the system is in a metastable condition, the competing forces operating on it are in a stable equilibrium. These forces maintain the system in the initially occupied energy well. There are two crucial variables for the loop [11, 12]: which indicates the magnetization obtained after the applied field has been removed. The second is coercive field, which is the field required to convert remanent magnetization to zero. A hysteresis loop's shape is dictated by the microscopic magnetic domain structure that corresponds to it [13]. The balance of several competing energy terms leads to the formation of magnetic domains. These competing energy terms include exchange energy, magneto-crystalline anisotropy, magneto static energy, which favours a null macroscopic average magnetic moment, and the interaction with external field, which favours magnetization along the direction of applied field. A continuous model created by Landau and Lifshitz forms the foundation of a generally recognised theoretical paradigm for ferromagnetism [14]. A universal energy functional, which incorporates the aforementioned energy factors, is used in their theory to theoretically account for magnetic phenomena on all pertinent length scales, including coarse domain structures, domain barriers, and their interactions. Due to the linkage with Maxwell's equations, the interaction of non-convexity and non-locality is what theoretically drives the generation of domain patterns [15, 16, 17].

In magnetic systems, there is usually a significant amount of structural disorder: polycrystalline grains, dislocations and lattice deformations, surface and interface roughness, compositional fluctuations, etc. Through exchange, anisotropy, and magneto static interactions, these sources of disorder are related to magnetism to produce an extraordinarily complex energy landscape with a large number of local minima that are almost equal in energy. The use of statistical approaches in the prediction of the magnetic properties of interest is therefore strongly recommended because domain structure and hysteresis are dominated by stochastic aspects [18, 19, 20].

REVIEW OF LITERATURE

It indicates a comparison of how similar a system's current configuration is to its former configuration. The system has a good memory if the average persistence of a certain microscopic configuration is high; otherwise, it has a bad memory [21]. Due to the complexity of such systems, a variety of factors influence memory, and as a result of their combined effects, memory is fundamentally a statistical quantity. Due to the widespread use of magnetic films in the modern information industry, memory research on magnetic films has the ability to not only provide a characteristic parameter for recording devices but also a method to increase the efficiency of magnetic storage. Madelung introduced the idea of macroscopically return point memory (RPM) in 1905, which may be the earliest magnetic memory concept [22, 23].

It focuses on the magnetic system's macroscopic memory, where magnetization is entirely controlled by applied fields. A rate independent main loop has perfect macroscopic RPM at every location of magnetization. Barkhausen noise revealed the existence of ferromagnetic domains, and since domains can now be seen using modern magnetic imaging techniques, it is now possible to study memory that exists on the length scale of the domain structure [24]. The topological domain configuration corresponding to a location on the hysteresis loop and the configuration corresponding to the same point after traversing cycles are compared to determine the microscopic RPM. Comparing domain configurations between conjugate points on a hysteresis loop is what microscopic conjugate point memory (CPM) does. Numerous elements, including anisotropy, interface roughness, and sample defects, might influence the microscopic RPM and CPM [25]. While some external characteristics, such as temperature and the applied magnetic field are controlled and can be used to study how the microscopic memory changes as a function of these variables, some are intrinsic and difficult to regulate. It was discovered that the RPM in the EB sample is high in the coercive area of the magnetization cycle and at the scale of the domain periodicity. Additionally, an odd spatial RPM oscillation was seen, suggesting a superstructure in the memory correlation. Despite the fact that exchange bias has been known about for more than 50 years, no one theory has been able to explain where it came from due to the fact that it can be found in a wide range of systems that have varying degrees of control over the structure, including cooling techniques, crystallinity, and growth methods. Dealing with interface structure is crucial for creating a complete theory because EB is linked to the exchange interaction formed at the FM-AFM interface [26, 27, 28, 29, 30].

DISCUSSION

Cellular and topological disorders are the two broad categories that comprise structural disorder. Small positional deviations from locations of perfect order, such as crystalline order, caused by flaws, thermal motion, etc. are referred to as cellular disorder. Topological disorder results from a group of things being randomly distributed without a particular order, like translational order. Topologically disordered ferromagnetic materials may exhibit complete spin order, supporting the notion of property-predefined order. Quasicrystals are disordered in terms of translational order, but they can also have some rotational symmetry that is prohibited by crystalline order in addition to long-range orientation arrange.

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REFERENCE

1. M. Cross and H. Greenside, *Pattern Formation and Dynamics in Nonequilibrium Systems* (Cambridge University Press, New York, 2009).
2. J. P. Gollub and J. S. Langer, *Rev. Mod. Phys.* 71, S396 (1999), URL <http://link.aps.org/doi/10.1103/RevModPhys.71.S396>.
3. M. C. Cross and P. C. Hohenberg, *Rev. Mod. Phys.* 65, 851 (1993), URL <http://link.aps.org/doi/10.1103/RevModPhys.65.851>.
4. B.-S. Han, D. Li, D.-J. Zheng, and Y. Zhou, *Phys. Rev. B* 66, 014433 (2002), URL <http://link.aps.org/doi/10.1103/PhysRevB.66.014433>.
5. Kreyssig, R. Prozorov, C. D. Dewhurst, P. C. Can_eld, R. W. McCallum, and A. I. Goldman, *Phys. Rev. Lett.* 102, 047204 (2009), URL <http://link.aps.org/doi/10.1103/PhysRevLett.102.047204>.
6. R. Rungswang, J. da Silva, C.-P. Wu, E. Sivaniah, A. Ionescu, C. H. W. Barnes, and N. J. Darton, *Phys. Rev. Lett.* 104, 255703 (2010), URL <http://link.aps.org/doi/10.1103/PhysRevLett.104.255703>.
7. J. Dickstein, S. Erramilli, R. E. Goldstein, D. P. Jackson, and S. A. Langer, *Science* 261, 1012 (1993), <http://www.sciencemag.org/content/261/5124/1012.full.pdf>, URL <http://www.sciencemag.org/content/261/5124/1012.abstract>.
8. <http://www.sciencemag.org/content/261/5124/1012.abstract>.
9. J. Richardi, D. Inger, and M. P. Pileni, *Phys. Rev. E* 66, 046306 (2002), URL <http://link.aps.org/doi/10.1103/PhysRevE.66.046306>.
10. G. Fosa, G. Dtrat, R. Badescu, and C. Gh., *J OPTOELECTRON ADV M* 6, 655 (2004).
11. G. Bertotti, *Hysterisis in Magnetism for Physicists, Materials Scientists, and Engineers* (Academic Press, Landon, 1998).

12. P. M. Ossi, *Disordered Materials: An introduction* (Springer-Verlag, Berlin, 2003). J. D. BERNAL, *Nature* 185, 68 (1960), URL <http://dx.doi.org/10.1038/185068a0>.
13. L. _Cervinka, *Journal of Non-Crystalline Solids* 106, 291 (1988), URL <http://www.sciencedirect.com/science/ARTICLE/pii/0022309388902773>.
14. P. H. Gaskell, M. C. Eckersley, A. C. Barnes, and P. Chieux, *Nature* 350, 675 (1991), URL <http://dx.doi.org/10.1038/350675a0>.
15. S. R. Elliott, *Nature* 354, 445 (1991), URL <http://dx.doi.org/10.1038/354445a0>.
16. J. M. Gibson and M. M. J. Treacy, *Phys. Rev. Lett.* 78, 1074 (1997), URL <http://link.aps.org/doi/10.1103/PhysRevLett.78.1074>.
17. P. S. Salmon, *Nat Mater* 1, 87 (2002), URL <http://dx.doi.org/10.1038/nmat737>.
18. J. D. Martin, S. J. Goettler, N. Fosse, and L. Iton, *Nature* 419, 381 (2002), URL <http://dx.doi.org/10.1038/nature01022>.
19. T. C. Hufnagel, *Nat Mater* 3, 666 (2004), URL <http://dx.doi.org/10.1038/nmat1227>.
20. H. W. Sheng, W. K. Luo, F. M. Alamgir, J. M. Bai, and E. Ma, *Nature* 439, 419 (2006), URL <http://dx.doi.org/10.1038/nature04421>.
21. Hirata, P. Guan, T. Fujita, Y. Hirotsu, A. Inoue, A. R. Yavari, T. Sakurai, and M. Chen, *Nat Mater* 10, 28 (2011), URL <http://dx.doi.org/10.1038/nmat2897>.
22. Fultz and J. Howe, *Transmission Electron Microscopy and Diractometry of Materials* (Springer-Verlag, Berlin, 2007).
23. M. M. J. Treacy, J. M. Gibson, L. Fan, D. J. Paterson, and I. McNulty, *Reports on Progress in Physics* 68, 2899 (2005), URL <http://stacks.iop.org/0034-4885/68/i=12/a=R06>.
24. L. FAN, D. PATERSON, I. McNULTY, M. M. J. TREACY, and J. M. GIBSON, *Journal of Microscopy* 225, 41 (2007), ISSN 1365-2818, URL <http://dx.doi.org/10.1111/j.1365-2818.2007.01714.x>.
25. Y. Z. Wu, C. Won, A. Scholl, A. Doran, H. W. Zhao, X. F. Jin, and Z. Q. Qiu, *Phys. Rev. Lett.* 93, 117205 (2004), URL <http://link.aps.org/doi/10.1103/PhysRevLett.93.117205>.
26. S. J. L. Billinge and M. F. Thorpe, *Local Structure from Di_raction* (Kluwer Academic Publishers, New York, 2002).
27. M. Seul and M. J. Sammon, *Phys. Rev. Lett.* 64, 1903 (1990), URL <http://link.aps.org/doi/10.1103/PhysRevLett.64.1903>.
28. P. Smith, J. F. Douglas, J. C. Meredith, E. J. Amis, and A. Karim, *Phys. Rev. Lett.* 87, 015503 (2001), URL <http://link.aps.org/doi/10.1103/PhysRevLett.87.015503>.
29. O. Portmann, A. Vaterlaus, and D. Pescia, *Nature* 422, 701 (2003), URL <http://dx.doi.org/10.1038/nature01538>.
30. D. G. Barci and D. A. Stariolo, *Phys. Rev. Lett.* 98, 200604 (2007), URL <http://link.aps.org/doi/10.1103/PhysRevLett.98.200604>.